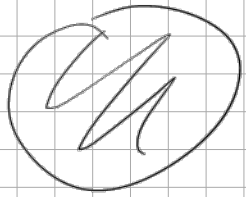


Status:

generation of
gravitational waves



$$T_{\mu\nu}(\vec{x}, t) = T_{\mu\nu}(\vec{x}) e^{-i\omega t} + \text{c.c.}$$

$$\begin{cases} \neq 0 & ; r = |\vec{x}| \leq R_0 \\ = 0 & ; r > R_0 \end{cases}$$

propagation of
gravitational wave

spherical wave in
far field $\approx$ planar wave



$$h_{\mu\nu}(x) = e_{\mu\nu} e^{-ik_\lambda x^\lambda} + \text{c.c.}$$

$$k_\lambda = \begin{pmatrix} \omega/c \\ k \vec{x}/r \end{pmatrix} \quad \omega = c k$$

angle dependent

$$e_{\mu\nu}(\vec{k}) = -\frac{4G}{c^4} \left(T_{\mu\nu}(\vec{k}) - \frac{T(\vec{k})}{2} \zeta_{\mu\nu} \right)$$

$$T_{\mu\nu}(\vec{k}) = \int d^3x T_{\mu\nu}(\vec{x}) e^{-i\vec{k} \cdot \vec{x}}$$

$$t_{\mu\nu}^{\text{grav}} = \frac{c^4}{16\pi G} k_\mu k_\nu \left(e^{\lambda\mu*} e^{\lambda\nu} - \frac{1}{2} |e^{\lambda\lambda}|^2 \right)$$

3.3.4 Radiated Power per Solid Angle:

local conservation law:

$$\partial^\mu t_{\mu\nu} = 0$$

$$r=0: \quad \frac{1}{c} \frac{\partial}{\partial t} t_{00} - \partial_i t_{i0} \Rightarrow \underbrace{\frac{\partial}{\partial t} \int d^3x t_{00}}_{\text{conserved energy}} = - \int d^3x \underbrace{\partial_i \left[\overbrace{t_{i0}}^{\text{energy current density}} \right]}_{\text{Gauss}} = 0$$

$$dP = -c \, t_{i0} \cdot \underbrace{d\Omega^i}_{\substack{\text{inf. solid angle}}} \\ = \frac{x^i}{r} r^2 \frac{d\Omega}{r^2}$$

$$= -c \cdot \frac{c^4}{16\pi G} \underbrace{(-x_i)}_{\substack{\text{w/c}}} \underbrace{x_0}_{\substack{\text{w/c}}} \left(e^{\lambda\kappa^*} e^{\lambda\kappa} - \frac{1}{2} |e^{\lambda\lambda}|^2 \right) \frac{x^i}{r} r^2 d\Omega \\ = -\frac{c^5}{16\pi G} \frac{1}{r} \left(e^{\lambda\kappa^*} e^{\lambda\kappa} - \frac{1}{2} |e^{\lambda\lambda}|^2 \right)$$

$$\frac{dP}{d\Omega} = \frac{c^5}{16\pi G} r^2 \left(e^{\lambda\kappa^*} e^{\lambda\kappa} - \frac{1}{2} |e^{\lambda\lambda}|^2 \right) \\ = \frac{c^5}{16\pi G} \frac{16 G^2}{c^8 r^2} \left\{ \left(\tau^{\lambda\kappa^*} - \frac{1}{2} \tau^{\lambda\lambda} \tau^{\kappa\kappa} \right) \left(\tau_{\lambda\kappa} - \frac{1}{2} \tau_{\lambda\lambda} \tau_{\kappa\kappa} \right) - \frac{1}{2} \left| \tau^{\lambda\lambda} - \frac{1}{2} \tau^{\lambda\lambda} \tau^{\lambda\lambda} \right|^2 \right\} \\ = \tau^{\lambda\kappa^*} \tau_{\lambda\kappa} - \frac{1}{2} \tau^{\lambda\lambda} \tau_{\lambda\lambda} - \frac{1}{2} \tau^{\kappa\kappa} \tau_{\kappa\kappa} + \frac{1}{2} \tau^{\lambda\lambda} \tau_{\lambda\lambda} \tau^{\kappa\kappa} \tau_{\kappa\kappa} - \frac{1}{2} \left| \tau^{\lambda\lambda} - \frac{1}{2} \tau^{\lambda\lambda} \tau^{\lambda\lambda} \right|^2$$

$$\frac{dP}{d\Omega} = \frac{G}{16\pi c^5} \left\{ \tau^{\lambda\kappa^*}(\vec{x}) \tau_{\lambda\kappa}(\vec{x}) - \frac{1}{2} \tau^{\lambda\lambda}(\vec{x}) \tau_{\lambda\lambda}(\vec{x}) \right\}$$

Ex 3.5 Reduction to Spatial Components:

$$\tau^{\mu\nu}(\vec{x}, t) = \int \frac{d^3k}{(2\pi)^3} \tau^{\mu\nu}(\vec{k}) e^{-i(\omega t - \vec{k} \cdot \vec{x})} \\ = k_{\lambda} x^{\lambda}$$

$$\partial_{\mu} \tau^{\mu\nu}(x^{\lambda}) = 0 \quad \xrightarrow{\text{weak fields}} \quad \partial_{\mu} \tau^{\mu\nu}(x^{\lambda}) = 0$$

Fourier $\rightarrow \quad \partial_\mu T^{\mu\nu}(\vec{k}) = 0$

$$v=0: \quad k_0 T^{00} = -k_i T^{i0} \quad (1)$$

$$v=i: \quad k_0 T^{i0} = -k_j T^{ij} \quad (2)$$

unit wave vector $\hat{k}_i = \frac{k_i}{k_0}$: contains angle dependencies (ϑ, φ)

$$(2) \quad T^{i0} = T^{0i} = -\hat{k}_i T^{ij}$$

$$(1) \quad T^{00} = -\hat{k}_i T^{i0} \quad \stackrel{(2)}{=} \hat{k}_i \hat{k}_j T^{ij}$$

$$1) \quad T^{\mu\nu*} T_{\mu\nu} = \sum_{\lambda\kappa} \hat{k}_{\lambda\kappa} T^{\mu\nu*} T^{\lambda\kappa}$$

$$= T^{00*} T^{00} - 2 T^{0i*} T^{0i} + T^{ij*} T^{ij}$$

$$= \left(\hat{k}_i \hat{k}_j \hat{k}_m \hat{k}_n - 2 \delta_{im} \hat{k}_j \hat{k}_n + \delta_{in} \delta_{jm} \right) T^{ij*} T^{mn}$$

$$) T^{ij*} T^{mn}$$

$$2) \quad T = T^{\lambda\lambda} = \sum_{\lambda\kappa} T^{\lambda\kappa} = T^{00} - T^{ii} = \left(\hat{k}_i \hat{k}_j - \delta_{ij} \right) T^{ij}$$

$$3) \quad T^* T = \left(\hat{k}_i \hat{k}_j \hat{k}_m \hat{k}_n - \delta_{ij} \hat{k}_m \hat{k}_n - \hat{k}_i \hat{k}_j \delta_{mn} + \delta_{ij} \delta_{mn} \right) T^{ij*} T^{mn}$$

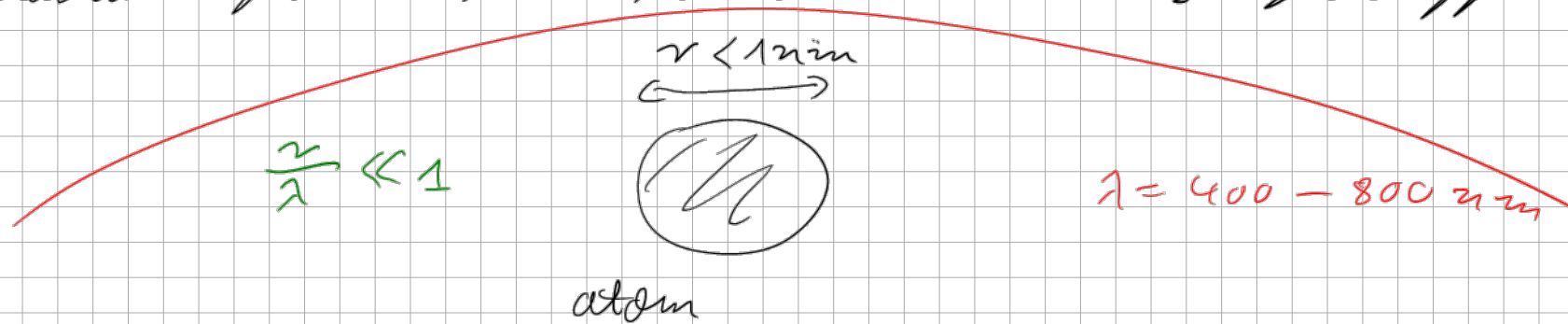
$$) T^{ij*} T^{mn}$$

$$\frac{dP}{d\Omega} = \frac{G \omega^2}{4\pi c^5} \underbrace{\mathcal{L}_{ij, mn}}_{\text{angle dependencies}} T^{ij*}(\vec{k}) T^{mn}(\vec{k})$$

$$A_{ij, mn} = \frac{1}{2} \hat{k}_i \hat{k}_j \hat{k}_m \hat{k}_n + \frac{1}{2} \delta_{ij} \hat{k}_m \hat{k}_n + \frac{1}{2} \delta_{mn} \hat{k}_i \hat{k}_j - 2 \delta_{im} \hat{k}_j \hat{k}_n \\ + \delta_{in} \delta_{jm} - \frac{1}{2} \delta_{ij} \delta_{mn}$$

3.3.6 Long-Wave Length Approximation:

In quantum optics this runs under the name of dipole approximation



Electromagnetic wave couples only to the electric dipole moment of the atom.

back to gravity:

$$|\vec{k} \cdot \vec{x}| \leq k r = \frac{2\pi}{\lambda} r \ll 1$$

$$T^{ij}(\vec{k}) = \int d^3x T^{ij}(\vec{x}) e^{-i\vec{k} \cdot \vec{x}} \quad \text{no longer } \vec{k}\text{-dependent}$$

$$\underbrace{e^{-i\vec{k} \cdot \vec{x}}}_{\text{small}} = 1 - i\vec{k} \cdot \vec{x} + \dots \approx 1$$

$$T^{ij}(\vec{x}) = \frac{1}{2} \int d^3x' x'^i x'^j \partial_{\mu} \partial_{\nu} T^{\mu\nu}(\vec{x}')$$

$$\partial_{\mu} T^{\mu\nu} = 0 \Rightarrow \begin{cases} \partial_0 T^{00} + \partial_j T^{j0} = 0 \\ \partial_0 T^{0i} + \partial_j T^{ji} = 0 \end{cases} \Rightarrow \partial_i \partial_j T^{ij} = \partial_0^2 T^{00} = -\frac{\omega^2}{c^2} T^{00}$$

$$\Rightarrow T^{ij}(\vec{x}) = -\frac{\omega^2}{2c^2} \int d^3x' x'^i x'^j T^{00}(\vec{x}')$$

$$T^{00}(\vec{x}) = \frac{\rho(\vec{x}) c^2}{1 - v^2/c^2} \quad (\text{special relativity})$$

$$v = \frac{r}{t} = \frac{1}{2\pi} \omega r = \frac{c}{2\pi} k r = \frac{c}{2\pi} \frac{2\pi}{\lambda} r = c \frac{r}{\lambda} \ll c$$

$\Rightarrow$ velocity contribution is negligible within long-wavelength approximation

$$\Rightarrow T^{ij}(\vec{x}) = -\frac{\omega^2}{2} Q^{ij}, \quad Q^{ij} = \int d^3x x'^i x'^j \rho(\vec{x}')$$

$$\Rightarrow \frac{dP}{d\Omega} = \frac{G \omega^6}{4\pi c^5} \Lambda_{ij, mn} Q^{ij} * Q^{mn}$$

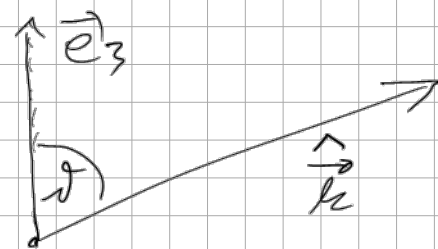
3.3.7 Radiated Power:

$$P = \int \frac{dP}{d\Omega} d\Omega = \frac{G \omega^6}{4\pi c^5} Q^{ij} * Q^{mn} \int d\Omega \Lambda_{ij, mn}$$

Question: moments

$$\int d\vec{r} \, 1 = 4\pi, \quad \int d\vec{r} \, \hat{k}_i \hat{k}_j = ?, \quad \int d\vec{r} \, \hat{k}_i \hat{k}_j \hat{k}_m \hat{k}_n = ?$$

$$\hat{k}_i = \frac{k_i}{k_0} = \begin{pmatrix} \sin\vartheta \cos\varphi \\ \sin\vartheta \sin\varphi \\ \cos\vartheta \end{pmatrix}$$



Method: generating function

$$g(\vec{x}) = \int d\vec{r} \, e^{-x_m x_n \hat{k}_m \hat{k}_n} = \int d\vec{r} \, e^{-(\vec{x} \cdot \hat{k})^2} = \int d\vec{r} \, e^{-r^2 \cos^2\vartheta}$$

$$= 2\pi \int_0^\pi d\vartheta \sin\vartheta \, e^{-r^2 \cos^2\vartheta} \quad \xrightarrow{u=\cos\vartheta} \quad 2\pi \cdot 2 \int_0^1 du \, e^{-r^2 u^2}$$

$$= \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} r^{2n} u^{2n} \quad \xrightarrow{x_i x_i x_j x_j} \quad = 4\pi \left\{ 1 - \frac{1}{3} r^2 + \frac{1}{10} r^4 + \dots \right\}$$

$$\Rightarrow g(\vec{x}) = 4\pi \sum_{n=0}^{\infty} \frac{(-1)^n}{(2n+1)n!} r^{2n} = 4\pi \left\{ 1 - \frac{1}{3} r^2 + \frac{1}{10} r^4 + \dots \right\}$$

$$g(\vec{x}=\vec{0}) = 4\pi \checkmark$$

$$-\frac{\partial^2 g(\vec{x})}{\partial x_i \partial x_j} \Big|_{\vec{x}=\vec{0}} = -\frac{\partial}{\partial x_i} \int d\vec{r} \, (-2 \hat{k}_i \hat{k}_n x_n) e^{-x_m x_n \hat{k}_m \hat{k}_n} \Big|_{\vec{x}=\vec{0}}$$

$$= \int d\vec{r} \, \{ 2 \hat{k}_i \hat{k}_j - 4 \hat{k}_i \hat{k}_j \hat{k}_m \hat{k}_n x_m x_n \} e^{-\dots} \Big|_{\vec{x}=\vec{0}}$$

$$= 2 \int d\vec{r} \, \hat{k}_i \hat{k}_j = \frac{4\pi}{3} \delta_{ij} \Rightarrow \int d\vec{r} \, \hat{k}_i \hat{k}_j = \frac{4\pi}{3} \delta_{ij}$$

$$\frac{\partial^4 f(\vec{x})}{\partial x_i \partial x_j \partial x_m \partial x_n} \Big|_{\vec{x}=\vec{0}} = \int d\Omega \left\{ \underbrace{(-2\vec{k}_i \vec{k}_j)(-2\vec{k}_m \vec{k}_n) + 4 \cdot 2 \vec{k}_i \vec{k}_j \vec{k}_m \vec{k}_n}_{= 12 \vec{k}_i \vec{k}_j \vec{k}_m \vec{k}_n} \right\}$$

$$= 8 \left\{ \delta_{ij} \delta_{mn} + \delta_{im} \delta_{jn} + \delta_{in} \delta_{jm} \right\}$$

$$\int d\Omega \vec{k}_i \vec{k}_j \vec{k}_m \vec{k}_n = \frac{4\pi}{15} \left(\begin{array}{c} \downarrow \end{array} \right)$$

$$\Rightarrow P = \frac{G \omega^6}{4\pi c^5} \frac{20}{15} Q^{ij*} Q^{mn} (-4 \delta_{ij} \delta_{mn} + 11 \delta_{im} \delta_{jn} + \delta_{in} \delta_{jm})$$

$$= \frac{2G \omega^6}{5 c^5} \left(Q^{is*} Q^{is} - \frac{1}{3} Q^{ii*} Q^{jj} \right)$$

compare with dipole radiation formula for electromagnetic waves

$$P = \frac{\omega^4}{12\pi \epsilon_0 c^4} \vec{p}_i^* \vec{p}_i$$

3.3.8 Amplitude:

We look for an amplitude of a planar gravitational wave, which has the same power P as the averaged one just calculated,

propagates in x^3 -direction:

$$(h^{\alpha}) = \frac{\omega}{c} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 1 \end{pmatrix}, \quad (e^{\alpha}_{\mu}) = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & e_{11} & e_{12} & 0 \\ 0 & e_{12} & -e_{11} & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$t^{\mu\nu}_{\text{grav}} = \frac{c^4}{8\pi G} h^{\mu\kappa} h^{\nu\lambda} (|e_{11}|^2 + |e_{12}|^2)$$

energy current densities per polarisation:

$$\underline{\Phi} = c t^{03} = \frac{c^3 \omega^2 h^2}{8\pi G} \quad (e_{11} = h, e_{12} = 0)$$

At distance D this amounts to the power

$$P = \underline{\Phi} 4\pi D^2 = \frac{c^3 \omega^2 h^2 D^2}{2G} \Rightarrow h = \sqrt{\frac{2GP}{c^3 \omega^2 D^2}}$$

3.3.9 Rotating Rigid Body:

rest frame: $(I'_{ij}) = \begin{pmatrix} I_1 & 0 & 0 \\ 0 & I_2 & 0 \\ 0 & 0 & I_3 \end{pmatrix}$

lab frame: $I(t) = \underset{L}{\alpha(t)} I' \alpha(t)^T = \text{const.} + Q_{ij} e^{-2i\omega t} + \text{c.c.}$

$$= \begin{pmatrix} \cos 2t & -\sin 2t & 0 \\ \sin 2t & \cos 2t & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad Q_{ij} = \frac{I_1 - I_2}{4} \begin{pmatrix} 1 & -i & 0 \\ -i & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

radiated power:

$$P = \frac{32 G \cancel{R^6}}{5 \cancel{c^5}} (\dot{I}_1 - \dot{I}_2)^2$$

→ small numbers

3.3.10 Rotating Bar

$$M = 500 \text{ t}, L = 20 \text{ m}, R = 30 \frac{1}{s}$$

$$I = \frac{m}{L}, \quad \dot{I}_1 = \frac{M}{12} L^2, \quad \dot{I}_2 \approx 0$$

$$\Rightarrow P = 3.6 \cdot 10^{-29} \text{ W} \quad \Rightarrow h = 2.2 \cdot 10^{-32}$$

⇒ no lab experiment yields observable gravitational waves

3.3.11 Double Star System:

→ ellipse $\approx$ circle

$$\dot{I}_1 = \frac{M_1 M_2}{M_1 + M_2} L^2, \quad L: \text{distance}$$

rotation frequency: gravitational force = centrifugal force

$$R^2 = G \frac{M_1 + M_2}{L^3}$$

$$\Rightarrow P = \frac{32 G^4}{5 c^5} \frac{M_1^2 M_2^2 (M_1 + M_2)}{L^5}$$



double star system: ϵ Boo

$$T = 0.268 \text{ d} \Rightarrow R = \frac{21}{T}$$

$$m_1 = 1.35 M_{\odot}, \quad m_2 = 0.68 M_{\odot}, \quad M_{\odot} = 1.55 \cdot 10^{30} \text{ kg}$$

$$\Rightarrow L = 1.5 \cdot 10^9 \text{ m} = 5.1 \text{ ls} = \frac{1}{100} \text{ AU}$$

astronomical unit
(distance sun - earth)

distance from earth: $D = 12 \text{ PC}$

nearest star to earth: Proxima Centauri = 12 PS

$$\Rightarrow L = 6.6 \cdot 10^{-21} \approx 10^{-20} \Rightarrow \underline{\text{detectable!}}$$